

# Method of eigenfunction expansion with homogeneous boundary conditions

We now seek to solve the following ~~equation~~  
nonhomogeneous heat equation (with homogeneous boundary conditions)

$$\left\{ \begin{array}{l} \text{(PDE)} \quad \frac{\partial v}{\partial t} = k \cdot \frac{\partial^2 v}{\partial x^2} + \bar{Q}(x, t) \quad \begin{array}{l} 0 \leq x \leq L \\ t > 0 \end{array} \\ \text{(*)} \left\{ \begin{array}{l} \text{(IC)} \quad v(0, t) = 0 \\ \quad \quad v(L, t) = 0 \\ \text{(IC)} \quad v(x, 0) = g(x) \end{array} \right. \end{array} \right.$$

Recall:

If  $\bar{Q}(x, t) = 0$ , then

$$\begin{aligned} v(x, t) &\equiv \sum_{n=1}^{\infty} B_n \cdot \sin\left(\frac{n\pi x}{L}\right) e^{-k \cdot \left(\frac{n\pi}{L}\right)^2 t} \\ &= \sum_{n=1}^{\infty} \underbrace{\left( B_n e^{-k \cdot \left(\frac{n\pi}{L}\right)^2 t} \right)}_{\equiv a_n(t)} \cdot \underbrace{\sin\left(\frac{n\pi x}{L}\right)}_{\phi_n(x)} \end{aligned}$$

where the related homogeneous boundary value problem

$$\frac{d^2 \phi}{dx^2} + \lambda \phi = 0$$

$$\phi(0) = 0$$

$$\phi(L) = 0$$

and

$$B_n = \frac{2}{L} \int_0^L g(x) \sin\left(\frac{n\pi x}{L}\right) dx$$

has the eigenvalues

$$\lambda_n = \left(\frac{n\pi}{L}\right)^2, \quad n = 1, 2, 3, \dots$$

with associated eigenfunctions

$$\phi_n(x) = \sin\left(\frac{n\pi x}{L}\right), \quad n = 1, 2, 3, \dots$$

In order to solve the nonhomogeneous heat equation (\*), the key idea is to write the unknown solution  $v(x, t)$  at every time  $t$  as an infinite sine series; in other words, to expand the unknown solution  $v(x, t)$  at every time  $t$  in terms of the eigenfunctions  $\phi_n(x)$  of the associated (homogeneous) boundary value problem

$$(**) v(x, t) = \sum_{n=1}^{\infty} a_n(t) \cdot \phi_n(x) \quad \text{for } t \geq 0$$

method of expansion  
eigenfunction

time-dependent  
(not equal to  $B_n e^{-k\left(\frac{n\pi}{L}\right)^2 t}$ )

$\leadsto$  For  $\phi_n(x) = \sin\left(\frac{n\pi x}{L}\right)$  this is just the Fourier sine series of  $v(x, t)$  at time  $t$ ; but we use this more general notation because the underlying idea works much more generally.

In (\*\*) the homogeneous boundary conditions are automatically satisfied and for the initial condition we need

$$g(x) = \sum_{n=1}^{\infty} a_n(0) \cdot \phi_n(x).$$

By the orthogonality of the eigenfunctions  $\phi_n(x)$  we can determine the coefficients

$$a_n(0) = \frac{\int_0^L g(x) \phi_n(x) dx}{\int_0^L \phi_n(x)^2 dx} \quad \left( = \frac{2}{L} \int_0^L g(x) \sin\left(\frac{n\pi x}{L}\right) dx \right)$$

To determine  $a_n(t)$  we plug the expansion (\*\*) for  $v(x,t)$  at every  $t \geq 0$  into the PDE for  $v$ :

$$\frac{\partial v}{\partial t} = \sum_{n=1}^{\infty} \frac{da_n}{dt}(t) \cdot \phi_n(x),$$

$$\begin{aligned} \frac{\partial^2 v}{\partial x^2} &= \sum_{n=1}^{\infty} a_n(t) \cdot \frac{d^2 \phi_n}{dx^2}(x) \\ &= - \sum_{n=1}^{\infty} a_n(t) \cdot \lambda_n \phi_n(x) \end{aligned}$$

Then our PDE

$$\frac{\partial v}{\partial t} - k \cdot \frac{\partial^2 v}{\partial x^2} = \bar{Q}(x,t)$$

reads

$$\sum_{n=1}^{\infty} \left( \frac{da_n(t)}{dt} + k \cdot \lambda_n a_n(t) \right) \phi_n(x) = \bar{Q}(x, t)$$

We can of course also expand the given (and thus known) source term  $\bar{Q}(x, t)$  in terms of the eigenfunctions  $\phi_n(x)$  (at every  $t \geq 0$ )

$$\bar{Q}(x, t) = \sum_{n=1}^{\infty} \bar{q}_n(t) \cdot \phi_n(x)$$

with

$$\bar{q}_n(t) = \frac{\int_0^L \bar{Q}(x, t) \phi_n(x) dx}{\int_0^L \phi_n(x)^2 dx} \left( = \frac{2}{L} \int_0^L \bar{Q}(x, t) \sin\left(\frac{n\pi x}{L}\right) dx \right)$$

Hence we have

$$\sum_{n=1}^{\infty} \left( \frac{da_n(t)}{dt} + k \cdot \lambda_n a_n(t) \right) \phi_n(x) = \sum_{n=1}^{\infty} \bar{q}_n(t) \phi_n(x)$$

and it follows that we must have for every  $n \in \mathbb{N}$

$$\frac{da_n(t)}{dt} + k \cdot \lambda_n a_n(t) = \bar{q}_n(t),$$

which is a 1<sup>st</sup> order linear nonhomogeneous ODE!

Multiplying it by the integrating factor  $e^{k\lambda_n t}$  we obtain

$$\underbrace{e^{k\lambda_n t} \left( \frac{da_n}{dt} + k\lambda_n a_n \right)} = e^{k\lambda_n t} \bar{q}_n(t)$$

$$= \frac{d}{dt} \left( e^{k\lambda_n t} a_n(t) \right)$$

Integrating from 0 to  $t$  yields

$$e^{k\lambda_n t} a_n(t) - a_n(0) = \int_0^t e^{k\lambda_n \tau} \bar{q}_n(\tau) d\tau$$

$$\Rightarrow a_n(t) = a_n(0) e^{-k\lambda_n t} + e^{-k\lambda_n t} \int_0^t e^{k\lambda_n \tau} \bar{q}_n(\tau) d\tau$$

With this choice of  $a_n(t)$  the solution to (\*) is then given by

$$v(x,t) = \sum_{n=1}^{\infty} a_n(t) \cdot \phi_n(x).$$

Note as a "sanity check" that if  $\bar{Q}(x,t) = 0$  we obtain

$$v(x,t) = \sum_{n=1}^{\infty} a_n(0) e^{-k\lambda_n t} \phi_n(x),$$

↖ Fourier sine coefficients of the initial condition  $f(x)$

which is what we'd had previously derived with the method of separation of variables.

Example:

Determine the solution to

$$\frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2} + e^{-t} \sin(3x) \quad \begin{array}{l} 0 \leq x \leq \pi \\ t > 0 \end{array}$$

$$u(0, t) = 0$$

$$u(\pi, t) = 1$$

$$u(x, 0) = f(x)$$

Step 1: We make the boundary conditions homogeneous by passing to the variable (displacement from equilibrium)

$$v(x, t) = u(x, t) - \frac{x}{\pi},$$

which has to satisfy

$$\frac{\partial v}{\partial t} = \frac{\partial^2 v}{\partial x^2} + e^{-t} \sin(3x)$$

$$v(0, t) = 0$$

$$v(\pi, t) = 0$$

$$v(x, 0) = f(x) - \frac{x}{\pi}$$

Step 2: Solve the nonhomogeneous heat equation for  $v(x, t)$  (with homogeneous BCs!)

The eigenfunctions are

$$\phi_n(x) = \sin\left(\frac{n\pi x}{\pi}\right) = \sin(nx), \quad n=1, 2, \dots$$

with eigenvalues

$$\lambda_n = \left(\frac{n\pi}{\pi}\right)^2 = n^2, \quad n=1, 2, \dots$$

So we write

$$v(x,t) = \sum_{n=1}^{\infty} a_n(t) \cdot \sin(nx).$$

Substituting into the PDE yields

$$\sum_{n=1}^{\infty} \left( \frac{da_n}{dt} + n^2 a_n \right) \sin(nx) = e^{-t} \sin(3x).$$

Thus,

$$\frac{da_n}{dt} + n^2 a_n = \begin{cases} 0 & , n \neq 3 \\ e^{-t} & , n = 3 \end{cases}$$

which has the solutions

$$\underline{\underline{n \neq 3:}} \quad a_n(t) = a_n(0) e^{-n^2 t}$$

$$\begin{aligned} \underline{\underline{n = 3:}} \quad a_3(t) &= a_3(0) e^{-9t} + e^{-9t} \int_0^t \underbrace{e^{+9\tau} \cdot e^{-\tau}}_{= e^{8\tau}} d\tau \\ &= a_3(0) e^{-9t} + \frac{1}{8} (e^{-t} + e^{-9t}) \end{aligned}$$

where

$$a_n(0) = \frac{2}{\pi} \int_0^{\pi} \left( f(x) - \frac{\pi}{\pi} \right) \sin(nx) dx.$$